

Geometric Maxwell-Lorentz : Discrete e⁻, p⁺

$$\{\hat{E}, \hat{B}\} = \{E, B\} / e \sim [l^{-2}]$$

$$\nabla \cdot \hat{E}(x) = 4\pi \sum_{j=1}^N s_j \delta(x - x_j) \quad s_j = \pm 1$$

$$\nabla \cdot \hat{B}(x) = 0$$

$$\begin{aligned} c\nabla \times \hat{E} &= -\dot{\hat{B}} \quad \text{photons} \\ c\nabla \times \hat{B} &= \dot{\hat{E}} + 4\pi \sum_{j=1}^N s_j v_j \delta(x - x_j) \end{aligned}$$

$$a_k = e^2 \frac{s_k}{m_k} \left[\hat{E} + \left(\frac{v_k}{c} \right) \times \hat{B} \right]$$

$$\mathcal{E} = \frac{e^2}{8\pi} (\hat{E}^2 + \hat{B}^2)$$

Fluidize : $N \rightarrow N/\varepsilon$ Constant : $\{E, B\}$
 $\varepsilon \rightarrow 0$

$$\{\hat{E}, \hat{B}\} \rightarrow \{\hat{E}, \hat{B}\} / \varepsilon$$

$$e \rightarrow \varepsilon e$$

$$m \rightarrow \varepsilon m$$

$$T \rightarrow \varepsilon T$$

$$r_e \rightarrow \varepsilon r_e \rightarrow 0$$

$$Q, \mathcal{E}$$

$$M, \bar{v}$$

$$nT$$

$$\text{MHD: } Q = 0$$

$$E = 0$$

$$\nabla \cdot J = 0$$

Fundamental Constants

$$\begin{array}{l} t \\ l \\ m \end{array} \left| \begin{array}{l} c \\ \\ m_e \end{array} \right.$$

$$e^2 = 1.44 \text{ eV} \cdot \text{nm}$$

$$hc = 1240. \text{ eV} \cdot \text{nm}$$

$$a_0 = \lambda_c^2 / r_e$$

$$\hat{\mu}_B = \lambda_c / 2$$

$$r_e \equiv \frac{e^2}{m_e c^2} = 2.82 \times 10^{-15} \text{ m}$$

$$\text{or } \lambda_c \equiv \frac{\hbar c}{m_e c^2} = 386. \times 10^{-15} \text{ m}$$

$$r_e / \lambda_c = 1 / 137.$$

$$\frac{G m_p^2}{e^2} = 0.81 \times 10^{-36}$$

Plasma Parameters :

Size scale r_e, n

Rate scale $\beta \equiv \bar{v} / c$

$$\bar{v}^2 \equiv T / m_e$$

$$\Omega_{ce} \equiv eB / mc = r_e \hat{B} c$$

$$\omega_p^2 \equiv 4\pi n e^2 / m = n r_e 4\pi c^2$$

$$\lambda_D^{-2} \equiv 4\pi n e^2 / T = n r_e \beta^{-2} 4\pi$$

$$" \beta_p^{-1} " \equiv \frac{B^2}{8\pi n T} = \frac{r_e \hat{B}^2 \beta^{-2}}{n 8\pi}$$

$$V_A^2 \equiv \frac{B^2}{4\pi n m_p} = \frac{r_e \hat{B}^2 c^2 m_e}{n 4\pi m_p}$$

Fluidize: $\rightarrow 0$

$$b \equiv e^2 / T = r_e \beta^{-2} \rightarrow 0$$

$$v_c \equiv n \bar{v} b^2 \ln \Lambda = n r_e^2 \beta^{-3} c \ln \Lambda \rightarrow 0$$

$$\rho_d \equiv \frac{m_e v_c}{e^2 n} = r_e \beta^{-3} \ln \Lambda / c \rightarrow 0$$

$$g \equiv \frac{1}{n \lambda_D^3} = [n r_e^3 \beta^{-6} (4\pi)^3]^{1/2} \rightarrow 0$$

$$\Gamma \propto g^{1/3} \text{ (correlation)}$$

Geometric Maxwell-Lorentz : Discrete e^- , p^+

In geometric units, $\{\hat{E}, \hat{B}\}$ have dimension length^{-2} .

The fundamental constant $e^2 = 1.44 \text{ eV}\cdot\text{nm}$ sets the EM energy scale;
or equivalently, $r_e = e^2/m_e c^2 = 2.8 \times 10^{-15} \text{ m}$ sets the (only) size scale .

The analogous quantum constant $hc = 1240. \text{ eV}\cdot\text{nm}$,
or equivalently the Compton wavelength λ_c ,
then determine the Bohr radius a_0 and the Bohr magneton μ_B .

Plasma parameters are then determined by
by the size scale set by r_e , density n , and geometric fields $\{\hat{E}, \hat{B}\}$,
and by the rate scale $\beta = \bar{v} / c$.

The fluid approximation “infinitely partitions” particles, keeping e/m constant;
but this makes $r_e = 0$, eliminating the only scale size in the equations, and
eliminating the physical basis for collisions, resistive dissipation, and correlation.
With separate e^- and p^+ fluids, dissipation can be re-inserted ad-hoc.

Magneto-Hydro Dynamics extends this mutilation of the Maxwell-Lorentz equations,
by eliminating all charge density, and all non-inductive electric fields,
retaining only a disembodied inductive (solenoidal, curling) current.